

Braids and Dynamics

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part 1
periodic motion

the taffy puller



Taffy is a type of candy.

Needs to be **pulled**: this aerates it and makes it lighter and chewier.

We can assign a **growth**: length multiplier per period.

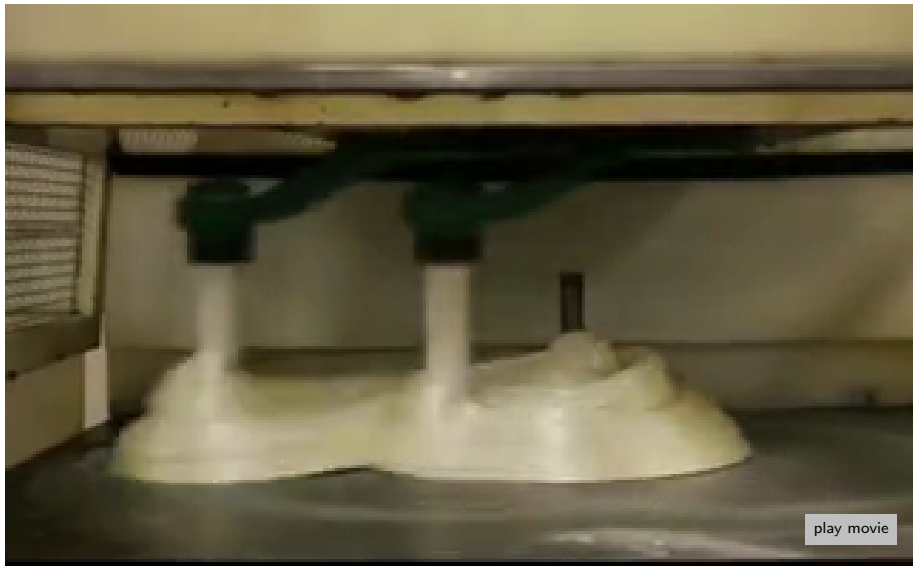
(Here $(1 + \sqrt{2})^2$... more later.)

[movie by M. D. Finn]

play movie



making candy cane



[*Wired*: This Is How You Craft 16,000 Candy Canes in a Day]

four-pronged taffy puller

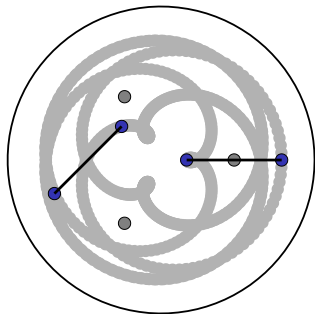


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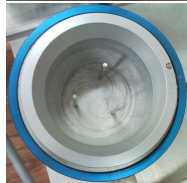
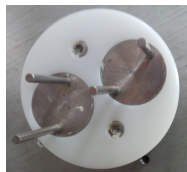
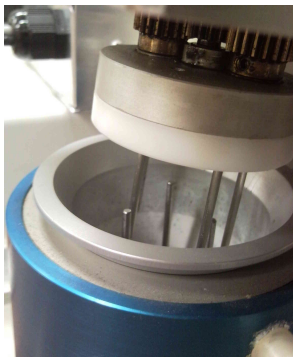
<http://www.youtube.com/watch?v=Y7t1HDSquVM>

[MacKay (2001); Halbert & Yorke (2014)]

Experimental device for kneading bread dough:



play movie



[Department of Food Science, University of Wisconsin. Photos by J-LT.]

For a *lot* more, see 'The Mathematics of Taffy Pullers,' Thiffeault, J.-L. (2018). *Math. Intelligencer*, **40** (1), 26–35. arXiv:1608.00152.

the mixograph as a braid



Encode the topological information as a sequence of **generators of the Artin braid group B_n** .

Equivalent to the 7-braid

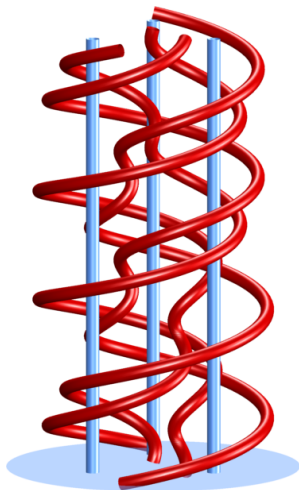
$$\sigma_3 \sigma_2 \sigma_3 \sigma_5 \sigma_6^{-1} \sigma_2 \sigma_3 \sigma_4 \sigma_3 \sigma_1^{-1} \sigma_2^{-1} \sigma_5$$

The **growth** is the largest root of

$$x^8 - 4x^7 - x^6 + 4x^4 - x^2 - 4x + 1$$

$$\simeq 4.186$$

Compare to taffy pullers: $\boxed{5.828}$





play movie



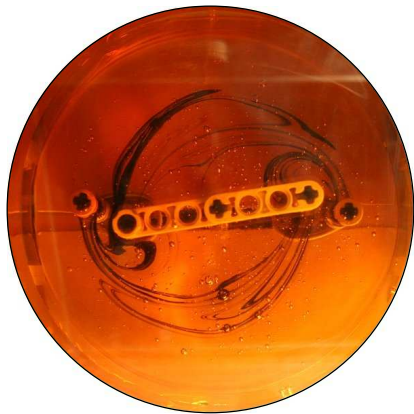
play movie

[Boyland, P. L., Aref, H., & Stremler, M. A. (2000). *J. Fluid Mech.* **403**, 277–304; Simulations by M. D. Finn, S. E. Tumas, and J-LT.]

4 + 1 rods



play movie



[Finn, M. D. & Thiffeault, J.-L. (2011). *SIAM Rev.* **53** (4), 723–743]





Periodic stirring protocols in two dimensions can be described by a **homeomorphism** $\varphi : \mathcal{S} \rightarrow \mathcal{S}$, where \mathcal{S} is a surface.

For instance, in a closed circular container,

- φ describes the mapping of fluid elements after one full period of stirring, obtained by solving the Stokes equation;
- \mathcal{S} is the **disc** with holes in it, corresponding to the stirring rods.

Goal: **Topological characterization of φ .**

[The theory extends to handlebodies, but not as relevant for applications. . .]

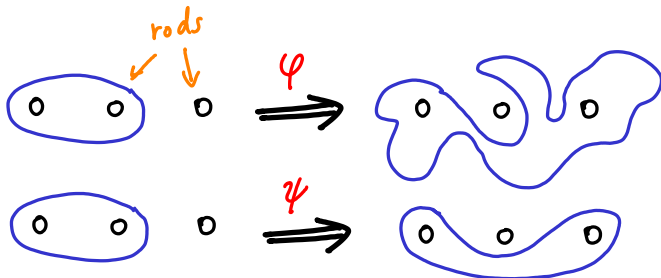


- 1 The Thurston–Nielsen classification theorem (**idealized φ**);
- 2 Handel's isotopy stability theorem (**link to real φ**);
- 3 Topological entropy (**quantitative measure of mixing**).

φ and ψ are **isotopic** if ψ can be continuously 'reached' from φ without moving the rods. Write $\varphi \simeq \psi$.

(Defines **isotopy classes**.)

Convenient to think of isotopy in terms of material loops. Isotopic maps act the same way on loops (up to continuous deformation).



(Loops will always mean **essential** loops.)



Theorem

φ is isotopic to a homeomorphism ψ , where ψ is in one of the following three categories:

finite-order for some integer $k > 0$, $\psi^k \simeq \text{identity}$;

reducible ψ leaves invariant a disjoint union of essential simple closed curves, called *reducing curves*;

pseudo-Anosov ψ leaves invariant a pair of transverse measured *singular foliations*, \mathcal{F}^u and \mathcal{F}^s , such that $\psi(\mathcal{F}^u, \mu^u) = (\mathcal{F}^u, \lambda \mu^u)$ and $\psi(\mathcal{F}^s, \mu^s) = (\mathcal{F}^s, \lambda^{-1} \mu^s)$, for *dilatation* $\lambda > 1$.

The three categories characterize the *isotopy class* of φ .

We want *pseudo-Anosov* for good mixing.

Handel's isotopy stability theorem



The TN classification tells us about a simpler map ψ , the **TN representative**. What about the original map φ ?

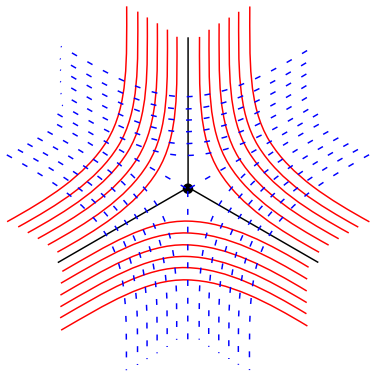
Theorem (Handel, 1985)

If ψ is pseudo-Anosov and isotopic to $\varphi : \mathcal{S} \rightarrow \mathcal{S}$, then there is a compact, φ -invariant set, $\mathcal{Y} \subset \mathcal{S}$, and a continuous, onto mapping $\alpha : \mathcal{Y} \rightarrow \mathcal{S}$, so that $\alpha\varphi = \psi\alpha$.

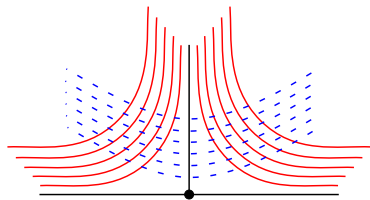
This is called a **semiconjugacy** (α not generally invertible).

Succinctly: the dynamics of the pseudo-Anosov map 'survive' isotopy, and so φ is at least as complicated as ψ . (In particular, **it has at least as much topological entropy.**)

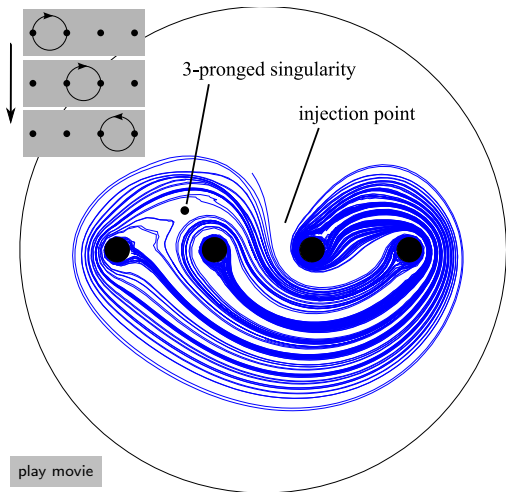
The 'pseudo' in pseudo-Anosov refers to the fact that the foliations can have a finite number of **pronged singularities**.



3-pronged singularity



Boundary singularity



play movie

- A four-rod stirring protocol;
- Material lines trace out leaves of the unstable foliation;
- Each rod has a 1-pronged singularity.
- One 3-pronged singularity in the bulk.
- One injection point (top): corresponds to boundary singularity;

[Thiffeault, J.-L., Finn, M. D., Gouillart, E., & Hall, T. (2008). *Chaos*, **18**, 033123]



- Consider a **motion of stirring elements**, such as rods.
- Determine if the motion is **isotopic to a pseudo-Anosov mapping**.
- **Compute** topological quantities, such as foliation, entropy, etc.
- **Analyze** and **optimize**.



[shameless plug for new book]



[Gouillart, E., Kuncio, N., Dauchot, O., Dubrulle, B., Roux, S., & Thiffeault, J.-L. (2007). *Phys. Rev. Lett.* **99**, 114501]

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ghost rods ('tiges fantômes')



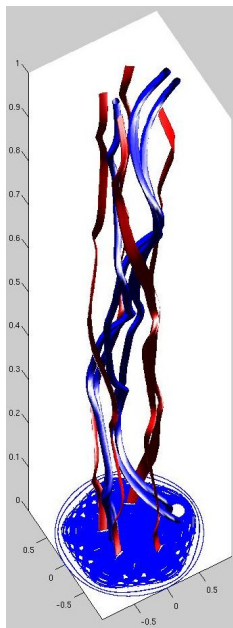
Topological analysis can be done on other objects than rods – for instance, **islands** or **unstable periodic orbits**.

We simply follow the islands and examine the braid they form, which gives us bounds on **topological entropy**.

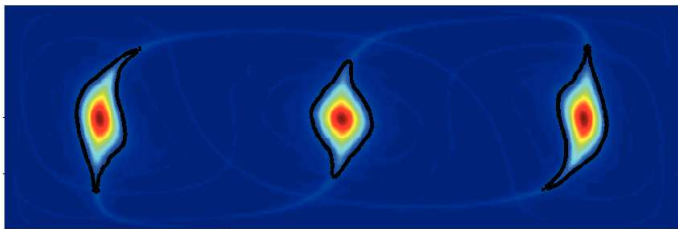
In this framework we call the islands **ghost rods**.

[Guillart, E., Finn, M. D., & Thiffeault, J.-L. (2006). *Phys. Rev. E*, **73**, 036311]

[implemented by Stremler & Chen (2007); Thiffeault *et al.* (2009); Binder (2010); Stremler *et al.* (2011)]



One of the best examples of ghost rods is from Stremler *et al.* (2011):



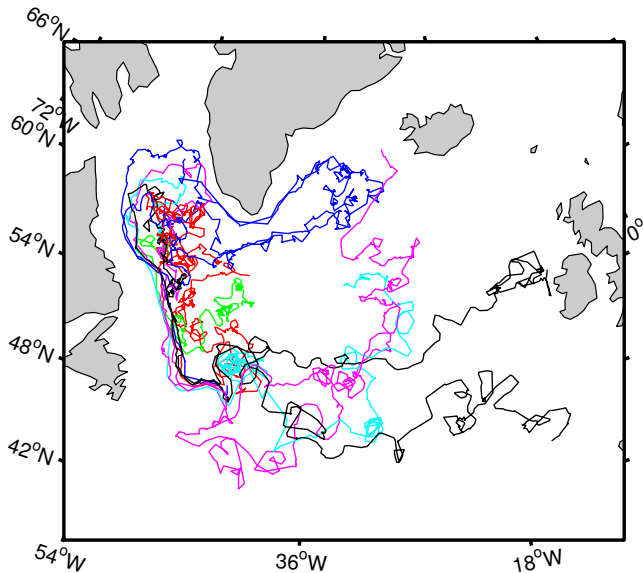
The islands are made to follow the $\sigma_2\sigma_1^{-1}$ stirring protocol by clever wall motions! (viscous Stokes flow)

[Stremler, M. A., Ross, S. D., Grover, P., & Kumar, P. (2011). *Phys. Rev. Lett.* **106**, 114101]

play movie

part 2
non-periodic motion

oceanic float trajectories





What can we measure?

- single-particle dispersion (not a good use of all data)
- correlation functions (useful)
- Lyapunov exponents (some luck needed!)

Another possibility:

Compute the **braid group generators** σ_i for the float trajectories (convert to a sequence of symbols), then look at how loops grow. Obtain a **topological entropy** for the motion (similar to Lyapunov exponent, or to the 'growth' of taffy pullers).



It is well-known that the entropy can be obtained by applying the motion of the punctures to a closed curve (loop) repeatedly, and measuring the growth of the length of the loop (Bowen, 1978).

The problem is twofold:

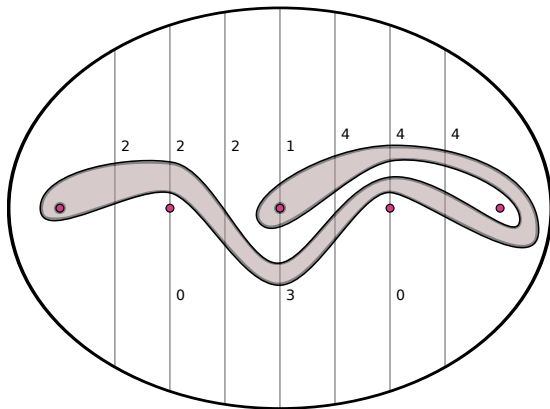
- 1 Need to keep track of the loop, since its length is growing exponentially;
- 2 Need a simple way of transforming the loop according to the motion of the punctures.

However, simple closed curves are easy objects to manipulate in 2D. Since they cannot self-intersect, we can describe them **topologically** with very few numbers.

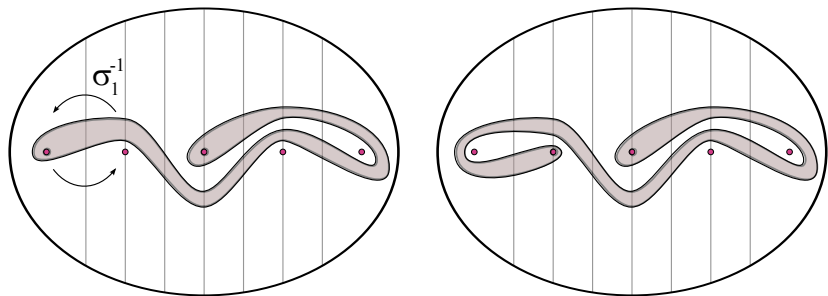
solution to problem 1: loop coordinates



What saves us is that a closed loop can be uniquely reconstructed from the number of intersections with a set of curves. For instance, the [Dybnikov coordinates](#) involve intersections with vertical lines:

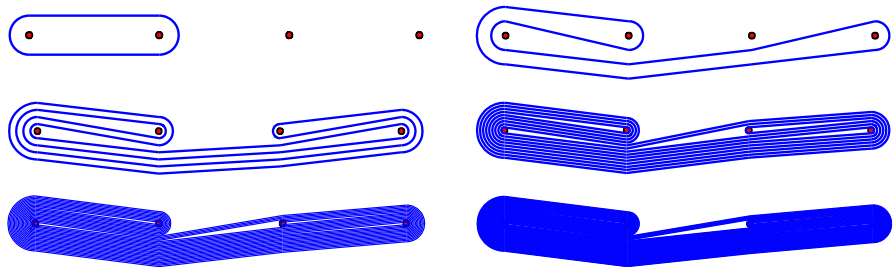


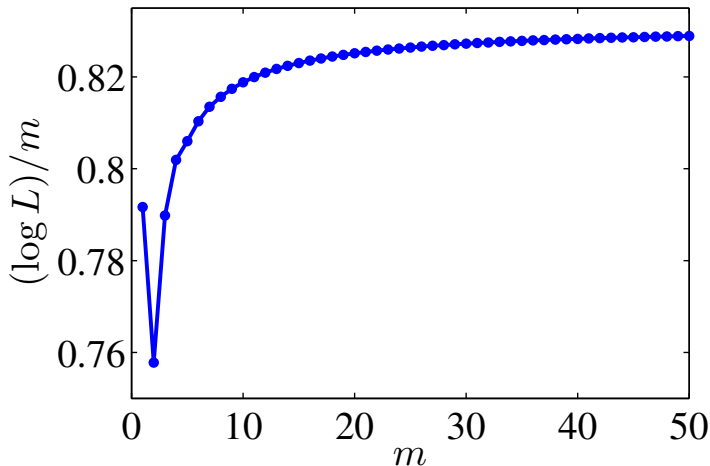
Moving the punctures according to a braid generator changes some crossing numbers:



There is an explicit formula for the change in the coordinates! [Dynnikov (2002); Moussafir (2006); Hall & Yurttas (2009); Thiffeault (2010)]

For a specific rod motion, say as given by the braid $\sigma_3^{-1}\sigma_2^{-1}\sigma_3^{-1}\sigma_2\sigma_1$, we can easily see the exponential growth of L and thus measure the entropy:

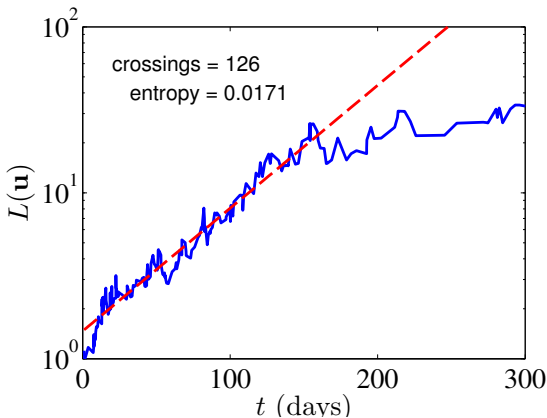




m is the number of times the braid acted on the initial loop.

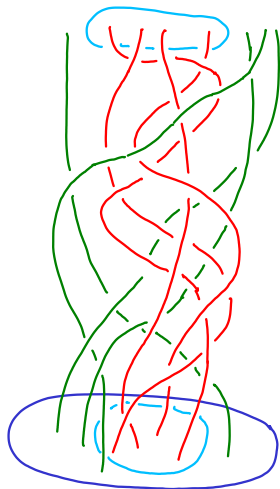
[Moussafir, J.-O. (2006). *Func. Anal. and Other Math.* 1 (1), 37–46]

10 floats from Davis' Labrador sea data:



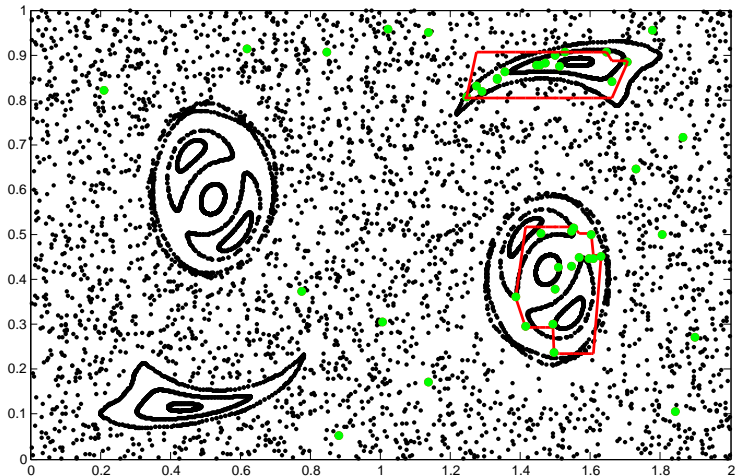
Floats have an entanglement time of about 50 days — timescale for horizontal stirring.

Source: WOCE subsurface float data assembly center (2004)



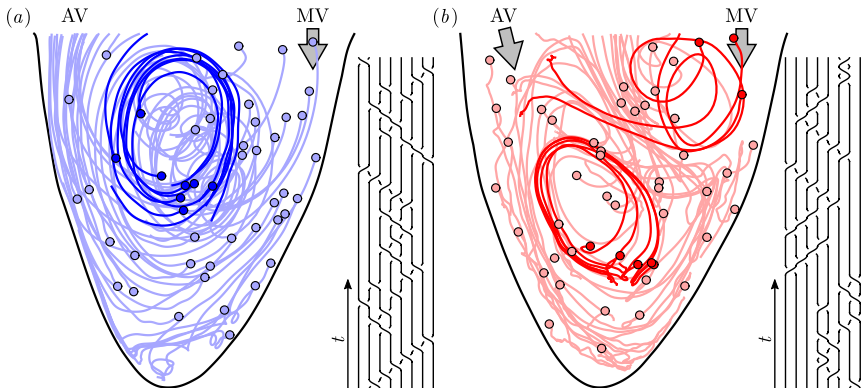
- There is a lot more information in the braid than just entropy;
- For instance: imagine there is an **isolated region** in the flow that does not interact with the rest, bounded by **Lagrangian coherent structures (LCS)**;
- Identify LCS and invariant regions from particle trajectory data by searching for curves that grow slowly or not at all.
- [see Haller, G. & Beron-Vera, F. J. (2012). *Physica D*, **241** (20), 1680–1702.]
- **Topological approach:** [Allshouse & Thiffeault (2012); Filippi *et al.* (2020); Yeung *et al.* (2020)].

double-gyre coherent structures



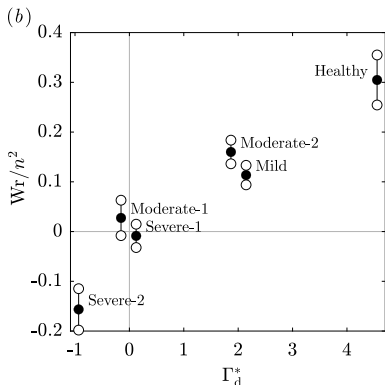
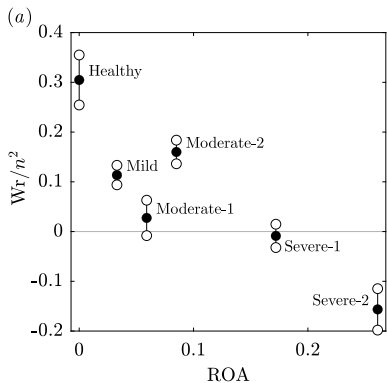
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[Allhouse, M. R. & Thiffeault, J.-L. (2012). *Physica D*, 241 (2), 95–105]



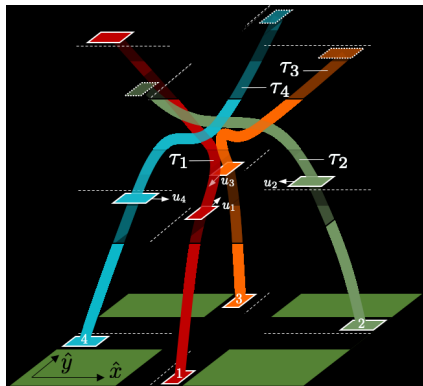
[Di Labbio, G., Thiffeault, J.-L., & Kadem, L. (2022). *Flow*, 2, E12]

braids in the heart (cont'd)



W_r writhe of braid
 ROA Regurgitant Orifice Area
 Γ_d^* circulation

[Di Labbio, G., Thiffeault, J.-L., & Kadem, L. (2022). *Flow*, 2, E12]



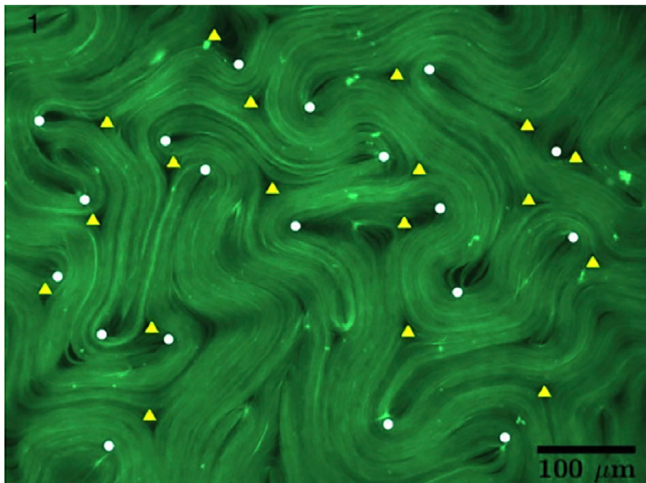
[Mavrogiannis, C., DeCastro, J., & Srinivasa, S. (2022). preprint]

This view of agent coordination is closely linked to Rob Ghrist's work on configuration space (U. Penn).

complexity of crowd movement



[Akpulat, M. & Ekinici, M. (2019). *Frontiers of Information Technology & Electronic Engineering*, 20 (6), 849–861]



[Smith, S. A. & Gong, R. (2022). *Frontiers in Physics*, **10**]



- We don't have solid theory for **aperodic** or **open** braids.
- **Computational methods** for isotopy class (random entanglements of trajectories – LCS method, see Allshouse & Thiffeault (2012); Filippi *et al.* (2020); Yeung *et al.* (2020)).
- **'Designing'** for topological chaos (see Stremler & Chen (2007)).
- Combine with **other measures**, e.g., **mix-norms** (Mathew *et al.*, 2005; Lin *et al.*, 2011; Thiffeault, 2012).
- Matlab toolbox — <https://github.com/jeanluct/braidlab>.
- **3D?!** (lots of missing theory; **E-Tec** approach shows promise [Roberts, E., Sindi, S., Smith, S. A., & Mitchell, K. A. (2019). *Chaos*, **29** (1), 013124]).

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