

Topological Stirring in Fluids

Jean-Luc Thiffeault

with

Matthew Finn and Emmanuelle Guillard

Department of Mathematics
Imperial College London

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The Taffy Puller

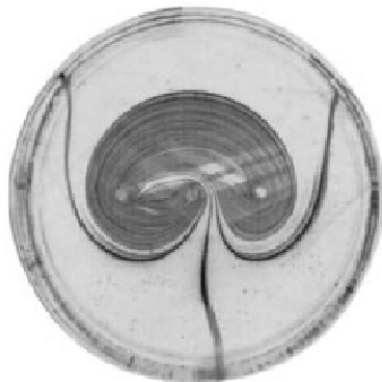


[movie 1]

The Four-pronged Taffy Puller



Experiment of Boyland, Aref, & Stremler

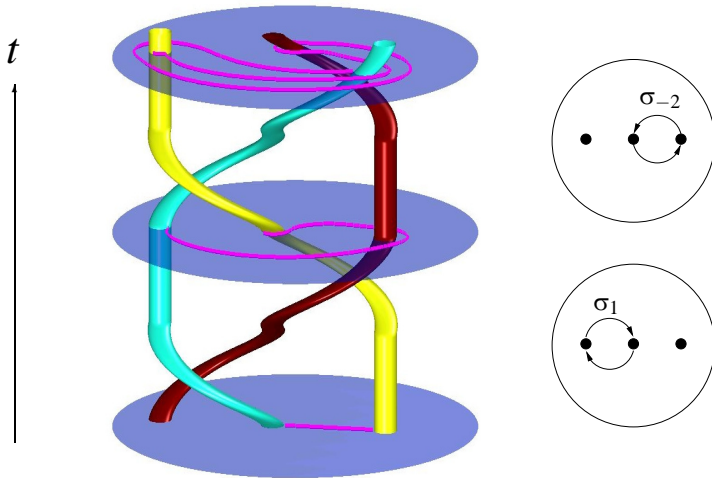


[P. L. Boyland, H. Aref, and M. A. Stremler, *J. Fluid Mech.* **403**, 277 (2000)]

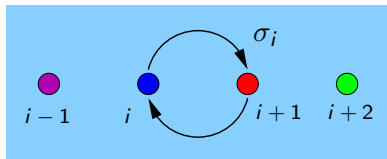
[P. L. Boyland, M. A. Stremler, and H. Aref, *Physica D* **175**, 69 (2003)]

[movie 2] [movie 3]

The Connection with Braids



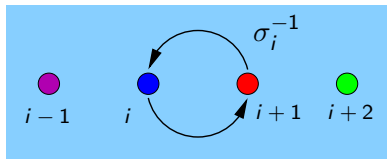
Generators of the n -Braid Group



A generator of Artin's braid group B_n on n strands, denoted

$$\sigma_i, \quad i = 1, \dots, n-1$$

is the clockwise interchange of the i th and $(i+1)$ th rod.



B_n is a **finitely-generated group**, with an infinite number of elements, called **words**.

These generators are used to characterise the topological motion of the rods.

Presentation of Artin's Braid Group

The generators obey the **presentation**

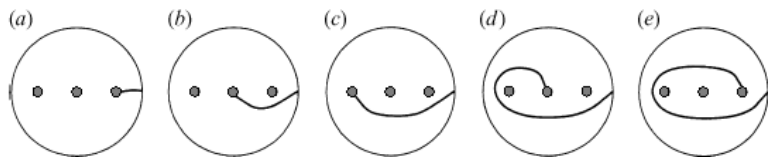
$$\sigma_{i+1} \sigma_i \sigma_{i+1} = \sigma_i \sigma_{i+1} \sigma_i$$

$$\sigma_i \sigma_j = \sigma_j \sigma_i, \quad |i-j| > 1$$

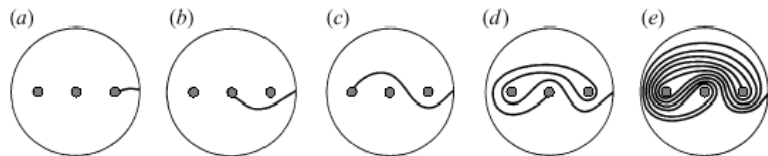
A presentation means that these are the **only rules** obeyed by the generators that are not the consequence of elementary group properties.

The Two BAS Stirring Protocols

$\sigma_1\sigma_2$ protocol



$\sigma_1^{-1}\sigma_2$ protocol

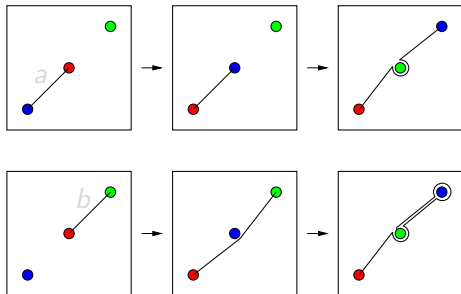
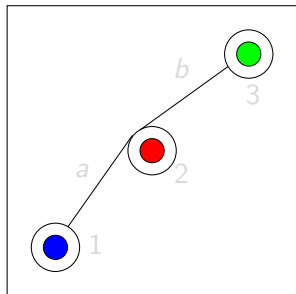


[P. L. Boyland, H. Aref, and M. A. Stremler, *J. Fluid Mech.* **403**, 277 (2000)]

Train-Tracks

What is the growth rate of an “elastic band” tied to the rods?

Train-tracks give the answer.



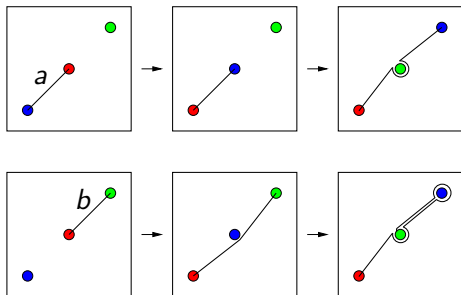
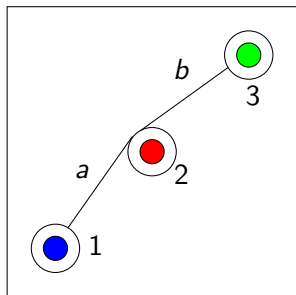
Elastic band has edges (letters) and infinitesimal loops (numbers).
As the rods are moved, the edges and loops are mapped as

$$a \mapsto a2b, \quad b \mapsto a2b3b, \quad 1 \mapsto 3, \quad 2 \mapsto 1, \quad 3 \mapsto 2.$$

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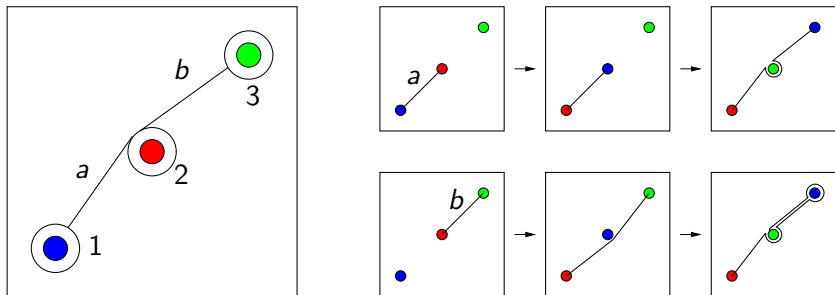
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The Evolution of Edges: Topological Entropy

The edges and loops are mapped according to

$$a \mapsto a2b, \quad b \mapsto a2b3b, \quad 1 \mapsto 3, \quad 2 \mapsto 1, \quad 3 \mapsto 2.$$

A crucial point is that edges are separated by loops: no cancellations can occur. A **transition matrix** can be formed:

$$M = \left[\begin{array}{cc|ccc} 1 & 1 & 0 & 0 & 0 \\ 1 & 2 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 1 & 0 \\ 1 & 1 & 0 & 0 & 1 \\ 0 & 1 & 1 & 0 & 0 \end{array} \right] \begin{array}{l} \leftarrow a \\ \leftarrow b \\ \leftarrow 1 \\ \leftarrow 2 \\ \leftarrow 3 \end{array}$$

The largest eigenvalue gives the asymptotic growth factor of the elastic is the **dilatation**, 2.6180. The logarithm of the dilatation is the **topological entropy** of the braid.

[M. Bestvina and M. Handel, *Topology* **34**, 109 (1995)]

The Difference between BAS's Two Protocols

- Practically speaking, the topological entropy of a braid is a lower bound on the **line-stretching exponent** of the flow!
- The first (**bad**) stirring protocol has zero topological entropy.
- The second (**good**) protocol has topological entropy $\log[(3 + \sqrt{5})/2] = 0.96 > 0$.
- **So for the second protocol the length of a line joining the rods grows exponentially!**
- That is, material lines have to stretch by at least a factor of 2.6180 each time we execute the protocol $\sigma_1^{-1}\sigma_2$.
- This is guaranteed to hold in some neighbourhood of the rods (**Thurston–Nielsen theorem**).

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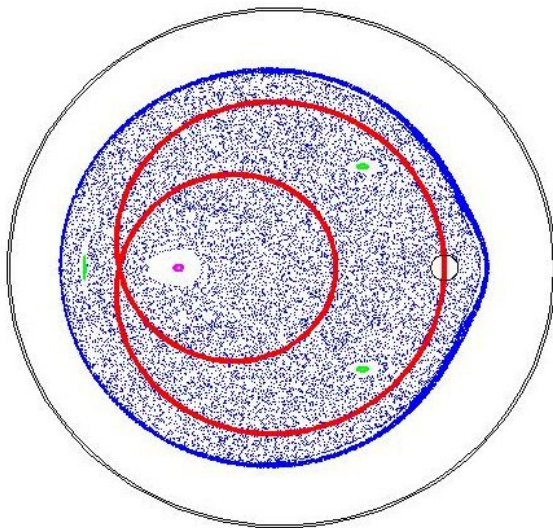
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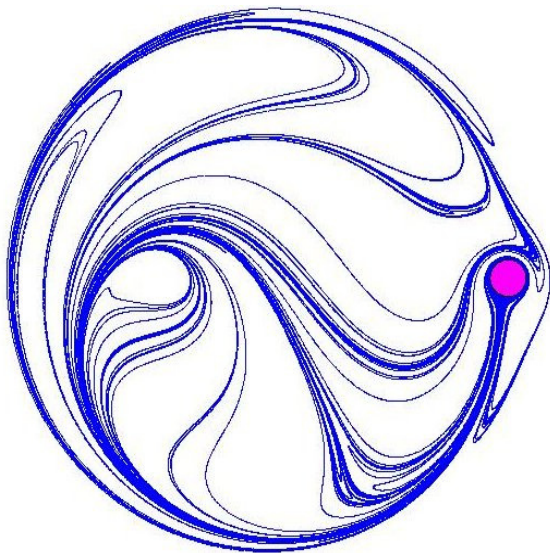
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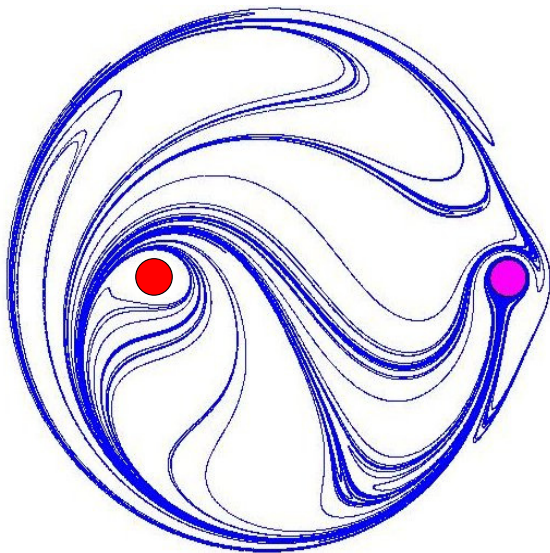
Poincaré Section



Stretching of Lines: A Ghostly Rod?

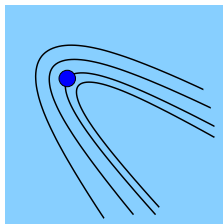
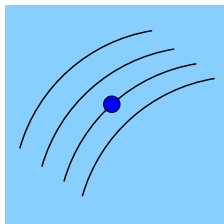


Stretching of Lines: A Ghostly Rod?



Particle Orbits are Topological Obstacles

Choose any fluid particle orbit (blue dot).



Material lines must bend around the orbit: **it acts just like a “rod”!**

[J-LT, *Phys. Rev. Lett.* **94**, 084502 (2005)]

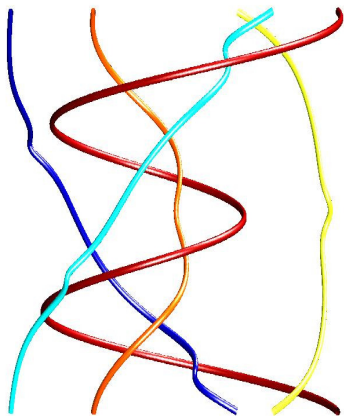
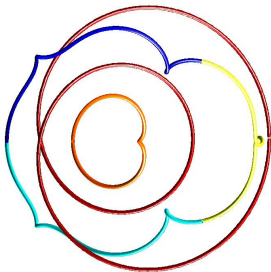
Today: focus on **periodic orbits**.

How do they braid around each other?

Motion of Islands

Make a braid from the motion of the rod and the **periodic islands**.

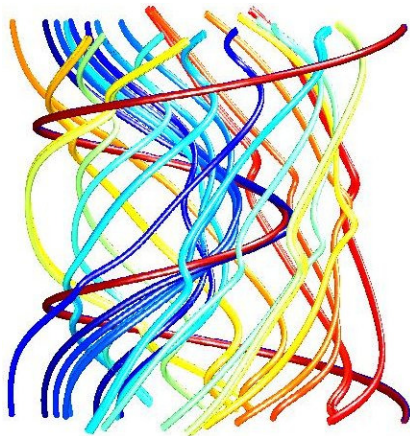
Most (74%) of the line-stretching is accounted for by this braid.



Motion of Islands and Unstable Periodic Orbits

Now we also include **unstable** periodic orbits as well as the stable ones (islands).

Almost all (99%) of the line-stretching is accounted for by this braid.



Optimisation

- The braid $\sigma_1\sigma_2^{-1}$ is **optimal** in B_3 . [Proved by d'Alessandro, Dahleh, and Mezić (1999)]
- The braid $\sigma_1\sigma_2^{-1}\sigma_3\sigma_2^{-1}$ is optimal in B_4 . [Conjecture]
- Both have dilatation $(1 + \sqrt{5})/2$ per generator, the **golden ratio**.
- In B_n , for $n > 4$, the golden ratio dilatation cannot be attained for an irreducible braid. [Conjecture]
- An artifact of the algebraic representation of the braid group: not very **physical**.
- This leads to a class of **silver ratio** mixers! (dilatation $1 + \sqrt{2}$)

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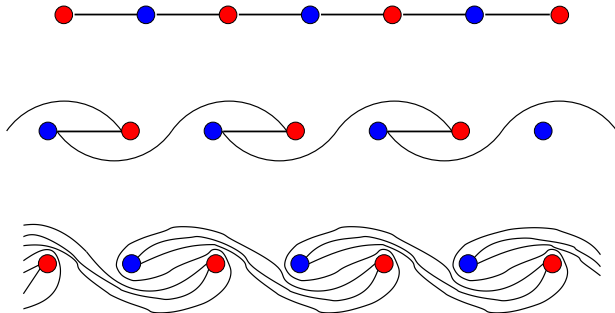
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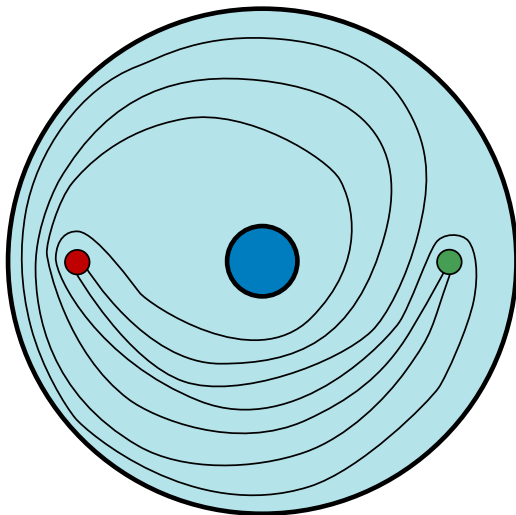
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Periodic Braid

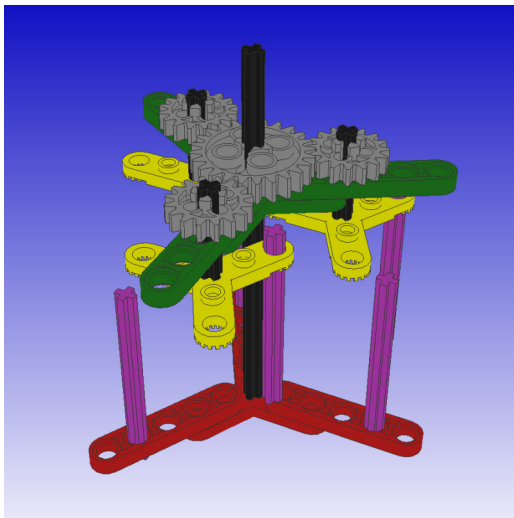


The dilatation of this braid is $1 + \sqrt{2}$, the [silver ratio](#).

Periodic Braid on Annulus



A Lego™ Implementation



[movie 4]

Conclusions

- Topological chaos involves moving **obstacles** in a 2D flow, which create nontrivial braids.
- A braid with positive topological entropy **guarantees** chaos in some region.
- Periodic orbits make great obstacles (in periodic flows), especially islands.
- This is a good way to “explain” the chaos in a flow — accounts for stretching of material lines.
- Other studies:
 - braids on the torus and sphere;
 - random braids;
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Recent Publications by Our Group

- T–G–F, “The Size of Ghost Rods,” in *Proceedings of the Workshop on Analysis and Control of Mixing* (Springer-Verlag, 2006, in press).
- G–T–F, “Topological Mixing with Ghost Rods,” *Physical Review E*, in press, 2006.
- F–T–G, “Topological Chaos in Spatially Periodic Mixers,” in submission, 2006.
- T, “Measuring Topological Chaos,” *Physical Review Letters* **94** (8), 084502, March 2005.

Preprints and slides available at www.ma.imperial.ac.uk/~jeanluc