# Modeling hagfish slime

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- Not the prettiest fish.
- An ancient fish: no teeth.
- Only known living animal that has a skull but not a vertebral column.
- 77 species, average 50 cm.
- Eats worms as well as dead fish, by burrowing into their carcass. They can feed through their own skin.



# sliming predators





#### play movie

[Zintzen, V., Roberts, C. D., Anderson, M. J., Stewart, A. L., Struthers, C. D., & Harvey, E. S. (2011). *Scientific Reports*, **1**, 131]

# knotting





youtube movie (see around 1 min mark)

# slime in the lab: a promising material





youtube movie

# so what's inside the slime?



- .002% thread skein
- .0015% mucin
- 99.996% seawater (!)



[Fudge, D. S., Levy, N., Chiu, S., & Gosline, J. M. (2005). J. Exp. Biol. 208, 4613-4625]



A skein consists of thread rolled into a ball.

Skeins are about 0.1 mm in size.

Thread length: about 15 cm!

The packing fraction is close to 1.



[Fernholm, B. (1981). Acta Zool. 62 (3), 137-145]

# what happens when the skeins unravel?

# The threads form a network, which gives the slime its properties.

The thread network determines the rheology of the slime.



[Fudge et al. (2005)]



# what happens when the skeins unravel?

Here the skein is stuck to a glass slide:



[experiment by Randy Ewoldt]



### similar to experiments with tape





The dynamics at the peeling points can get very complicated and can even lead to triboluminescence.

[Figures from Cortet, P.-P., Ciccotti, M., & Vanel, L. (2007). J. Stat. Mech. 2007, P03005 and Camara, C. G., Escobar, J. V., Hird, J. R., & Putterman, S. J. (2008). Nature, 455, 1089–1092; see also Maugis & Barquins (1988); De et al. (2004).]

Work-energy theorem of Hong & Yue (1995):

$$\dot{U} = (T - F_0(V))V$$

- $\dot{U}$  is the total change in the energy of the system;
- *T* is the force drawing out the thread;
- $F_0(V)$  is a velocity-dependent peeling force.

Neglect changes to the elastic energy of the tape ( $\dot{U} = 0$ ):

$$T=F_0(V)$$





A simple model for the peeling force is

$$F_0(V) = \alpha V^m, \qquad 0 \le m \le 1$$

which we solve for the peeling velocity:

$$V = (T/\alpha)^{1/m}$$

The total length L(t) of thread drawn out thus satisfies

$$\dot{L} = (T/\alpha)^{1/m}$$



Relate R, the skein radius, and L using mass conservation:

$$\frac{d}{dt}\left(\frac{4}{3}\pi\eta R^3 + \pi r^2 L\right) = 0 \quad \Longrightarrow \quad \dot{L} = -4\eta R^2 \dot{R}/r^2,$$

where r is the thread radius and  $\eta \leq 1$  is the packing fraction of thread into the spherical skein.

Under constant tension T, we can easily solve for the depletion time

$$t_{\rm dep} = \frac{4\eta R_0^3}{3r^2} \left( T/\alpha \right)^{1/m}$$

to run out of thread, given an initial skein radius  $R_0$ .



So far there is no fluid. Assume the skein is immersed a simple 1D flow

$$\mathbf{u}(x,y,t) = u(x,t)\,\hat{\mathbf{x}}$$

at the origin. The thread remains straight and aligned with the horizontal.

Resistive force theory for a straight filament then says

$$8\pi\mu\delta\left(\mathbf{x}_{t}-\mathbf{u}\right)=2T_{s}\mathbf{x}_{s}$$
$$2T_{ss}=-8\pi\mu\delta\mathbf{x}_{s}\cdot\mathbf{u}_{s}$$

where  $\mathbf{x}(s, t)$  is the Eulerian position of a thread segment as a function of the Lagrangian label s (arc length).

 $[\mu = \text{viscosity}, \ \delta = \text{slenderness parameter}]$ 





- The Lagrangian arc length parameter values s ∈ [0, L<sub>0</sub>] correspond to the 'initial' piece of thread.
- The thread added by unraveling is  $s \in (L_0, L(t)]$ .

The tension at s = L(t) — the end unraveling from the skein — is equal to the drag force on a sphere of radius R:

$$T = 6\pi\mu R \mathbf{x}_s \cdot (\mathbf{u} - \mathbf{x}_t), \qquad s = L(t).$$

The other end is free:

$$T=0, \qquad s=0.$$

The Eulerian position of a thread element is

$$x = X(t) - L(t) + s$$

where X(t) is the position of the skein.

Hence, using the tension in our peeling law:

$$(\dot{L})^m = 6\pi\mu R\alpha^{-1}(\dot{L} - \dot{X} + u(X, t))$$

This is not closed: we need to find a separate equation for X. We do this by solving the equations of resistive force theory.



Skipping some algebraic details, we eventually obtain the system of differential equations

$$(\dot{L})^{m} = 6\pi\mu\alpha^{-1}R\,\bar{u}(X,L,t)$$
$$\dot{X} = \dot{L} + u(X,t) - \bar{u}(X,L,t)$$

where

$$\bar{u}(X,L,t) = \frac{1}{L + (3R/2\delta)} \int_{X-L}^{X} \{u(X,t) - u(x,t)\} \, \mathrm{d}x$$

Note that R and L are related by mass conservation.





For an extensional flow

$$u(x,t) = \lambda(t)x$$

we can reduce the system to one ODE:

$$(\dot{L})^m = 3\pi\mu\alpha^{-1}\lambda RL^2/(L + (3R/2\delta))$$

where again R = R(L), and also  $\delta = \delta(L)$  (slenderness parameter).

[Recall:  $\delta = -1/\log(\varepsilon^2 e)$  with  $\varepsilon = r/L$  the slenderness ratio.]

This nonlinear ODE cannot be solved analytically, except in some asymptotic limits for special choices of m.

### numerical solution





Numerical solution for  $r = 1 \,\mu\text{m}$ ,  $R_0 = 50 \,\mu\text{m}$ ,  $L_0 = 2R_0$ ,  $\mu = 1.5 \times 10^{-3} \,\text{Pas}$ , m = 1/3,  $\alpha = 8 \times 10^{-4} \,\text{N} \,(\text{m/s})^{-1/3}$ ,  $\eta = 1$ ,  $\lambda = 1 \,\text{s}^{-1}$ .

# numerical solution: discussion



- The depletion time is about  $10^3$  s.
- Problem: this is 3 orders of magnitude too long!
- Possible issues:
  - No real idea what peeling force parameter values to use. (Here used educated guess based on tape.) Changing these can radically alter the results.
  - Maybe adjust the drag force if the filament doesn't remain straight: go beyond resistive force to full slender-body theory.
  - Are the mucins important? Experiments suggest so but their role is unclear. They might catalyze the peeling somehow, or stick to the filament and increase the drag force.
  - Proper rheological experiments needed (Randy Ewoldt)!
  - Use a more 'mixing' flow, closer to turbulence.
- Future research: network created by threads.

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