

Nonuniform mixing

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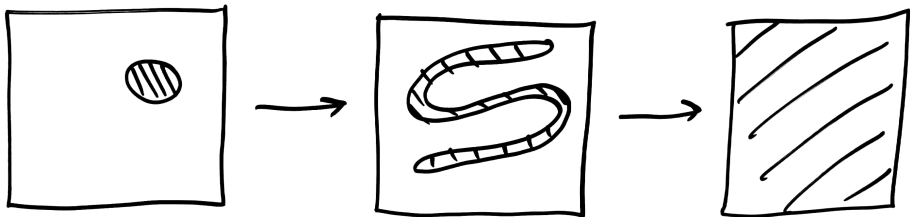
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The usual scenario in mixing is that we want to homogenize some initial distribution of **particles** or **dye**.



This will happen naturally via **molecular diffusion**, but is greatly accelerated by stirring.

See for instance Welander, P. (1955). *Tellus*, **7** (2), 141–156.



The **advection-diffusion equation** governs the evolution of a passive scalar concentration $\theta(\mathbf{x}, t)$:

$$\partial_t \theta + \mathbf{u} \cdot \nabla \theta = D \nabla^2 \theta, \quad \nabla \cdot \mathbf{u} = 0,$$

where $\mathbf{u}(\mathbf{x}, t)$ is a divergence-free **velocity field**, and D is the **diffusivity**.

With **no-flux boundary conditions**

$$(\mathbf{u} \theta - D \nabla \theta) \cdot \hat{\mathbf{n}} = 0$$

at the boundary $\partial\Omega$ of the domain Ω , the integral $\int_{\Omega} \theta \, dV$ is conserved.



How do we know that the concentration will **eventually mix**? A few integration by parts and use of boundary conditions give

$$\frac{d}{dt} \int_{\Omega} \theta^2 dV = -2D \int_{\Omega} |\nabla \theta|^2 dV \leq 0.$$

The decay of **variance** (L^2 norm) is monotonic: it can never increase. It can only stop decreasing if θ is uniform in space ($\nabla \theta \equiv 0$).

This bound underpins the usefulness of variance as a measure of mixing.



The **relaxation to a uniform state** requires this uniform state to be a steady solution of the advection–diffusion equation.

Perhaps surprisingly, **this is not always the case!**

A uniform state is a steady solution of the advection-diffusion equation only if $\nabla \cdot \mathbf{u} = 0$ (which we take as given), as well as

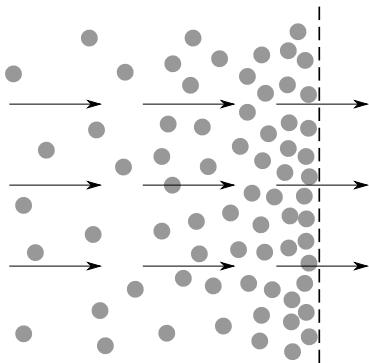
$$\mathbf{u} \cdot \hat{\mathbf{n}} = 0 \quad \text{on the boundary } \partial\Omega.$$

For $\mathbf{u} \cdot \hat{\mathbf{n}} \neq 0$, the uniform state $\theta = \text{const.}$ solves the advection–diffusion equation, but does not satisfy the boundary conditions. The equilibrium state is **nonuniform**.

An example: particle filter



The simplest example of this is a **filter**: $\mathbf{u} \cdot \hat{\mathbf{n}} \neq 0$ at the boundary, since fluid can cross the filter, but the **particles cannot**.



The equilibrium state is then **nonuniform**: particles tend to accumulate at suction regions on the boundary.



With $\mathbf{u} \cdot \hat{\mathbf{n}} \neq 0$ on the boundary, the evolution of variance is now given by

$$\frac{d}{dt} \int_{\Omega} \theta^2 dV = \int_{\partial\Omega} \theta^2 \mathbf{u} \cdot \hat{\mathbf{n}} dS - 2D \int_{\Omega} |\nabla\theta|^2 dV.$$

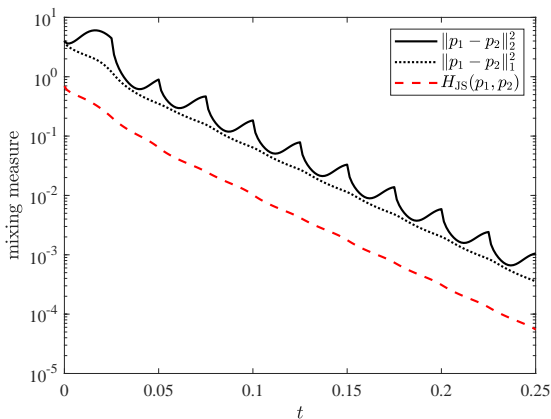
Note the boundary term on the right is **not sign-definite**. Hence variance no longer has to decrease monotonically. It can exhibit transient growth.

Of course **variance must ultimately decay**, which we know from other considerations. But the above equation does not show that, and suggests that variance **can be poorly-behaved if used as a measure of mixing**.

Relaxation to equilibrium: Example



A simple example is a constant flow $U\hat{x}$ on the interval $[0, 1]$. We apply no-flux boundary conditions and periodically reverse the direction of the flow ('breathing').



Notice that variance (solid line) shows significant oscillations.

A better measure of mixing in the nonuniform case is the f -divergence:

$$H_f[p_1, p_2] := \int_{\Omega} p_2 f(p_1/p_2) dV.$$

Here p_1 and p_2 are two **normalized probability densities**, and f is a **convex function** with $f(1) = 0$, $f'' \geq 0$.

For example we can choose

$$f(u) = u \log u$$

which gives the **Kullback–Leibler divergence** or **relative entropy**.

H_f measures the ‘distance’ (divergence) between p_1 and p_2 . We set $p_1 = \theta(\mathbf{x}, t)$, and p_2 to the steady solution.



The reason f -divergence is a nice measure of mixing is that

$$\frac{d}{dt} H_f[p_1, p_2] = -D \int_{\Omega} p_2 f''(p_1/p_2) |\nabla(p_1/p_2)|^2 dV \leq 0$$

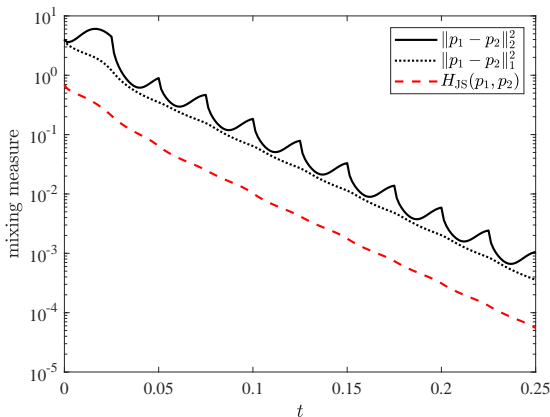
for general no-flux boundary conditions, that is, even if $\mathbf{u} \cdot \hat{\mathbf{n}} \neq 0$. The relaxation of f -divergence is thus **always monotonic**.

This is essentially an **H -theorem** from statistical physics. The novelty here is that in those applications the boundary conditions are not important, since quantities such as momentum vanish at infinity. In the fluid-dynamical context it is **precisely the no-flux boundary conditions** that give this monotonic evolution of H_f .

Relaxation of f -divergence: Example



Return to the earlier periodic flow example: the dashed red line is the f -divergence. Notice how nice and monotonic it is compared to variance (solid).



The dotted line is the L^1 norm $\int_{\Omega} |p_1 - p_2| dV$, and is **also monotonic!**



The previous plot suggests that the L^1 norm

$$\int_{\Omega} |p_1 - p_2| dV = \int_{\Omega} |\theta| dV$$

is also monotonic in time.

Indeed, this follows from

$$\frac{d}{dt} \int_{\Omega} |\theta| dV = -2D \int_{\{\theta=0\}} |\nabla\theta| dS \leq 0$$

where the integral on the right is taken over the **zero level set** of $\theta(\cdot, t)$.

This again holds **even in the nonuniform case**, but it is less useful mathematically. In practice, it suggests that L^1 is **a more reliable measure of mixing than variance** for nonuniform mixing.



- Mixing is usually regarded as the **relaxation to a uniform state**.
- The **concentration variance** (L^2 norm) is often taken as a **convenient measure**, since it relaxes **monotonically** to a uniform state.
- However, in some cases the ultimate state is **not uniform**.
- For example: **suction boundary conditions**, or divergent flows (not discussed).
- In those nonuniform cases **variance is less reliable**, since it can exhibit oscillations: **it is not constrained to decay monotonically**.
- Better measures of mixing in the nonuniform case are the entropy-like quantities called **f -divergence**, or the **L^1 norm**.
- See Thiffeault, J.-L. (2021). *Physical Review Fluids*, **6** (9), 090501.



- Risken, H. (1996). *The Fokker–Planck Equation: Methods of Solution and Applications*. Berlin: Springer, second edition.
- Thiffeault, J.-L. (2021). *Physical Review Fluids*, **6** (9), 090501.
- Welander, P. (1955). *Tellus*, **7** (2), 141–156.