

# Supplemental Material for “Energy Decompositions for Moist Boussinesq and Anelastic Equations with Phase Changes”

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## Energy Principle for the Kessler Scheme

In this section, we will derive the energy principle given in section 4a of the main text. First, we obtain an expression for the material derivative of  $\Pi$ . Using the fact that

$$\frac{D\Pi}{Dt} = \frac{\partial\Pi}{\partial b_u^{\text{tot}}} \frac{Db_u^{\text{tot}}}{Dt} + \frac{\partial\Pi}{\partial b_s^{\text{tot}}} \frac{Db_s^{\text{tot}}}{Dt} + \frac{\partial\Pi}{\partial z} \frac{Dz}{Dt}, \quad (1)$$

we have

$$\begin{aligned} \frac{D}{Dt}\Pi(b_u^{\text{tot}}, b_s^{\text{tot}}, z) &= -[(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] \frac{Dz}{Dt} \\ &\quad - \int_a^z \frac{\partial}{\partial b_u^{\text{tot}}} [(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] dz' \frac{Db_u^{\text{tot}}}{Dt} \\ &\quad - \int_a^z \frac{\partial}{\partial b_s^{\text{tot}}} [(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] dz' \frac{Db_s^{\text{tot}}}{Dt}. \end{aligned} \quad (2)$$

The second and third terms on the right hand side of (2) can be simplified further:

$$\begin{aligned} & - \int_a^z \frac{\partial}{\partial b_u^{\text{tot}}} [(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] dz' \\ &= - \int_a^z H_u + (b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u(z') - \tilde{b}_s(z'))) \frac{\partial}{\partial b_u^{\text{tot}}} H_u dz' \\ &= - \int_a^z H_u dz' - \int_a^z (b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u(z') - \tilde{b}_s(z'))) \frac{\partial}{\partial b_u^{\text{tot}}} H_u dz' \\ &= - \int_a^z H_u dz' \end{aligned} \quad (3)$$

where we have used the fact that

$$\int_a^z (b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u(z') - \tilde{b}_s(z'))) \frac{\partial}{\partial b_u^{\text{tot}}} H_u dz' = 0$$

A similar procedure yields

$$- \int_a^z \frac{\partial}{\partial b_s^{\text{tot}}} [(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] dz' = - \int_a^z H_s dz'. \quad (4)$$

We can integrate the terms  $\int_a^z H_u dz'$  and  $\int_a^z H_s dz'$  that appear (3) and (4) exactly. Using an integration by parts we have:

$$\begin{aligned}
& - \int_a^z H_u dz' = -zH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u(z') - \tilde{b}_s(z'))) \Big|_a^z \\
& - \int_a^z z' (N_u^2 - N_s^2)(z') \delta(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z')) dz' \\
& = -zH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z)) + aH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \\
& - \int_{-\infty}^{\infty} [H(z' - a) - H(z' - z)] z' (N_u^2 - N_s^2)(z') \delta(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z')) dz' \\
& = -zH_u + aH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) - z_r H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) + z_r H_u.
\end{aligned} \tag{5}$$

A similar procedure results in

$$\begin{aligned}
- \int_a^z H_s dz' & = -(z - a) + zH_u - aH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \\
& + z_r H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) - z_r H_u.
\end{aligned} \tag{6}$$

We have thus shown, so far, that

$$\begin{aligned}
- \int_a^z \frac{\partial}{\partial b_u^{\text{tot}}} [(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] dz' & = -zH_u + aH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \\
& - z_r H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) + z_r H_u
\end{aligned} \tag{7}$$

and

$$\begin{aligned}
- \int_a^z \frac{\partial}{\partial b_s^{\text{tot}}} [(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] dz' & = -(z - a) + zH_u - aH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \\
& + z_r H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) - z_r H_u,
\end{aligned} \tag{8}$$

in which case,

$$\begin{aligned}
& - \int_a^z \frac{\partial}{\partial b_u^{\text{tot}}} [(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] dz' \frac{Db_u^{\text{tot}}}{Dt} \\
& = \left[ -zH_u + aH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) - z_r H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) + z_r H_u \right] \\
& \left( \frac{g}{\tilde{\rho}(z)} \left( R_{vd} - \frac{L_v}{c_p \theta_0} \right) \frac{\partial}{\partial z} (\tilde{\rho}(z) V_T q_r) - g \left( R_{vd} - \frac{L_v}{c_p \theta_0} + 1 \right) S_r \right),
\end{aligned} \tag{9}$$

and

$$\begin{aligned}
& - \int_a^z \frac{\partial}{\partial b_s^{\text{tot}}} [(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s] dz' \frac{Db_s^{\text{tot}}}{Dt} = \\
& - \left[ -(z - a) + zH_u - aH(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) + z_r H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) - z_r H_u \right] \\
& \frac{g}{\tilde{\rho}(z)} \frac{\partial}{\partial z} (\tilde{\rho}(z) V_T q_r),
\end{aligned} \tag{10}$$

where  $S_r = A_r + C_r - E_r$ .

Combining (10) and (9), and a great deal of simplifying we obtain the material derivative of  $\Pi$ :

$$\begin{aligned}
\frac{D}{Dt}\Pi(b_u^{\text{tot}}, b_s^{\text{tot}}, z) &= -[(b_u^{\text{tot}} - \tilde{b}_u)H_u + (b_s^{\text{tot}} - \tilde{b}_s)H_s]w + (z - a)\frac{gV_T}{\tilde{\rho}(z)}\frac{\partial}{\partial z}(\tilde{\rho}(z)q_r) \\
&\quad -g(z - z_r)\left(R_{vd} + \frac{L_v}{c_p\theta_0} + 1\right)\left[\frac{1}{\tilde{\rho}(z)}\frac{\partial}{\partial z}(\tilde{\rho}(z)V_Tq_r) - S_r\right]H_u \\
-g(z_r - a)\left(R_{vd} - \frac{L_v}{c_p\theta_0} + 1\right)\left[\frac{1}{\tilde{\rho}(z)}\frac{\partial}{\partial z}(\tilde{\rho}(z)V_Tq_r) - S_r\right] &H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a))
\end{aligned} \tag{11}$$

Using the momentum equations, it is straightforward to verify that

$$\tilde{\rho}(z)\frac{D}{Dt}\left(\frac{1}{2}|\mathbf{u}|^2\right) = -\tilde{\rho}(z)\nabla\phi + \tilde{\rho}(z)w(b_uH_u + b_sH_s). \tag{12}$$

Adding  $\tilde{\rho}(z)D\Pi/Dt$  to (12), and using incompressibility, we obtain the desired equation:

$$\begin{aligned}
\frac{\partial}{\partial t}\left(\tilde{\rho}(z)\left(\frac{|\mathbf{u}|^2}{2} + \Pi\right)\right) + \nabla \cdot \left[\tilde{\rho}(z)\mathbf{u}\left(\frac{|\mathbf{u}|^2}{2} + \Pi + \phi\right)\right] &= g(z - a)\frac{\partial}{\partial z}(V_T\tilde{\rho}(z)q_r) \\
-g(z - z_r)\left(R_{vd} - \frac{L_v}{c_p\theta_0} + 1\right)\left[\frac{\partial}{\partial z}(V_T\tilde{\rho}(z)q_r) - \tilde{\rho}(z)S_r\right] &H_u \\
-g(z_r - a)\left(R_{vd} - \frac{L_v}{c_p\theta_0} + 1\right)\left[\frac{\partial}{\partial z}(V_T\tilde{\rho}(z)q_r) - \tilde{\rho}(z)S_r\right] &H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)).
\end{aligned} \tag{13}$$

## Integrating $\Pi$ for the Anelastic Equations

In this appendix, we provide the details of the decomposition of  $\Pi$  described in section 3b of the main text. We integrate the term  $\int_a^z (b_u^{\text{tot}} - \tilde{b}_u)H_u dz'$  by parts:

$$\begin{aligned}
-\int_a^z (b_u^{\text{tot}} - \tilde{b}_u(z'))H_u dz' &= -H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z'))\int_a^{z'} (b_u^{\text{tot}} - \tilde{b}_u(s)) ds \Big|_{z'=a}^{z'=z} \\
-\int_a^z (N_u^2(z') - N_s^2(z'))\delta(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z')) &\int_a^{z'} (\tilde{b}_u^{\text{tot}} - \tilde{b}_u(s)) ds dz'.
\end{aligned} \tag{14}$$

The first term in (14) is straightforward to evaluate:

$$\begin{aligned}
-H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z'))\int_a^{z'} (b_u^{\text{tot}} - \tilde{b}_u(s)) ds \Big|_{z'=a}^{z'=z} \\
= -H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s))\int_a^z (b_u^{\text{tot}} - \tilde{b}_u(s)) ds
\end{aligned} \tag{15}$$

For the second term in (14), we multiply the integrand by the characteristic function of  $[a, z]$  and integrate over the entire real line. Doing this ensures that the zero of the delta function's argument lies in the region of integration.

$$\begin{aligned}
-\int_a^z (N_u^2(z') - N_s^2(z'))\delta(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z'))\int_a^{z'} (\tilde{b}_u^{\text{tot}} - \tilde{b}_u(s)) ds dz' &= \\
-\int_{-\infty}^{\infty} \left[ [H(z' - a) - H(z - z')] (N_u^2(z') - N_s^2(z')) \right. \\
\left. \delta(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z')) \int_a^{z'} (\tilde{b}_u^{\text{tot}} - \tilde{b}_u(s)) ds \right] dz' &
\end{aligned} \tag{16}$$

Now make the substitution  $u = b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z')$ . Then  $du = -(N_u^2(z') - N_s^2(z')) dz'$ ,

$$\lim_{z' \rightarrow \infty} b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z') = -\infty, \quad (17)$$

$$\lim_{z' \rightarrow -\infty} b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z') = \infty, \quad (18)$$

and the integral on the right hand side of (16) becomes

$$\begin{aligned} & - \int_{-\infty}^{\infty} \left[ [H(z' - a) - H(z - z')] (N_u^2(z') - N_s^2(z')) \right. \\ & \left. \delta(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z')) \int_a^{z'} (\tilde{b}_u^{\text{tot}} - \tilde{b}_u(s)) ds \right] dz' = \\ & \int_{-\infty}^{\infty} \left( H((\tilde{b}_u - \tilde{b}_s)^{-1}(u') - a) - H((\tilde{b}_u - \tilde{b}_s)^{-1}(u') - z) \right) \delta(u) \int_a^{(\tilde{b}_u - \tilde{b}_s)^{-1}(u')} (b_u^{\text{tot}} - \tilde{b}_u(s)) ds du, \end{aligned} \quad (19)$$

where  $u' = b_u^{\text{tot}} - b_s^{\text{tot}} - u$ . The integral on the right hand side can be evaluated by setting  $u = 0$ , yielding

$$\begin{aligned} & \int_{-\infty}^{\infty} \left[ \left( H((\tilde{b}_u - \tilde{b}_s)^{-1}(u') - a) - H((\tilde{b}_u - \tilde{b}_s)^{-1}(u') - z) \right) \delta(u) \right. \\ & \left. \int_a^{(\tilde{b}_u - \tilde{b}_s)^{-1}(u')} (b_u^{\text{tot}} - \tilde{b}_u(s)) ds \right] du = \\ & - [H((\tilde{b}_u - \tilde{b}_s)^{-1}(b_u^{\text{tot}} - b_s^{\text{tot}}) - a) - H((\tilde{b}_u - \tilde{b}_s)^{-1}(b_u^{\text{tot}} - b_s^{\text{tot}}) - z)] \times \\ & \int_a^{(\tilde{b}_u - \tilde{b}_s)^{-1}(b_u^{\text{tot}} - b_s^{\text{tot}})} (b_u^{\text{tot}} - \tilde{b}_u(s)) ds = \\ & - [H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) - H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z))] \int_a^{z_r} (b_u^{\text{tot}} - \tilde{b}_u(s)) ds \end{aligned} \quad (20)$$

where  $z_r = (\tilde{b}_u - \tilde{b}_s)^{-1}(b_u^{\text{tot}} - b_s^{\text{tot}})$ , and we have used the facts that

$$H((\tilde{b}_u - \tilde{b}_s)^{-1}(b_u^{\text{tot}} - b_s^{\text{tot}}) - a) = H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \quad (21)$$

and

$$H((\tilde{b}_u - \tilde{b}_s)^{-1}(b_u^{\text{tot}} - b_s^{\text{tot}}) - z) = H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z)). \quad (22)$$

To see why equations (21) and (22) are true, note that  $\tilde{b}_u - \tilde{b}_s$ , and  $(\tilde{b}_u - \tilde{b}_s)^{-1}$  are both monotone increasing functions. A consequence of this is that  $(\tilde{b}_u - \tilde{b}_s)^{-1}(b_u^{\text{tot}} - b_s^{\text{tot}}) - z \geq 0$  if and only if  $b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z) \geq 0$ , which ensures that  $H((\tilde{b}_u - \tilde{b}_s)^{-1}(b_u^{\text{tot}} - b_s^{\text{tot}}) - z) = H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(z))$ . Therefore,

$$\begin{aligned} - \int_a^z (b_u^{\text{tot}} - \tilde{b}_u(z')) H_u dz' &= -H_u \int_a^z (b_u^{\text{tot}} - \tilde{b}_u(z')) dz' + H_u \int_a^{z_r} (b_u^{\text{tot}} - \tilde{b}_u(z')) dz' \\ & - H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \int_a^{z_r} (b_u^{\text{tot}} - \tilde{b}_u(z')) dz' \end{aligned} \quad (23)$$

A similar procedure yields

$$\begin{aligned} - \int_a^z (b_s^{\text{tot}} - \tilde{b}_s(z')) H_s dz' &= -H_s \int_a^z (b_s^{\text{tot}} - \tilde{b}_s(z')) dz' - H_u \int_a^{z_r} (b_s^{\text{tot}} - \tilde{b}_s(z')) dz' \\ & + H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \int_a^{z_r} (b_s^{\text{tot}} - \tilde{b}_s(z')) dz' \end{aligned} \quad (24)$$

Adding these two expressions and making further simplifications we obtain

$$\begin{aligned}
\Pi &= -H_u \int_a^z (b_u^{\text{tot}} - \tilde{b}_u(z')) dz' - H_s \int_a^z (b_s^{\text{tot}} - \tilde{b}_s(z')) dz' \\
&+ H_u \int_a^{z_r} (b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u(z') - \tilde{b}_s(z'))) dz' \\
&- H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \int_a^{z_r} (b_u^{\text{tot}} - \tilde{b}_u(z')) dz' \\
&+ H(b_u^{\text{tot}} - b_s^{\text{tot}} - (\tilde{b}_u - \tilde{b}_s)(a)) \int_a^{z_r} (b_s^{\text{tot}} - \tilde{b}_s(z')) dz'.
\end{aligned} \tag{25}$$