Biomixing

when organisms stir their environment

Jean-Luc Thiffeault\textsuperscript{1}  Steve Childress\textsuperscript{2}  Zhi George Lin\textsuperscript{3}

\textsuperscript{1}Department of Mathematics
University of Wisconsin – Madison

\textsuperscript{2}Courant Institute of Mathematical Sciences
New York University

\textsuperscript{3}Department of Mathematics
Zhejiang University

Complex Systems Seminar
University of Michigan, Ann Arbor, 4 December 2012

Supported by NSF grants DMS-0806821 and DMS-1109315
Biomixing

A controversial proposition:

- There are many regions of the ocean that are relatively quiescent, especially in the depths (1 hairdryer/ km$^3$);  
- Yet mixing occurs: nutrients eventually get dredged up to the surface somehow;  
- What if organisms swimming through the ocean made a significant contribution to this?  
- There could be a local impact, especially with respect to feeding and schooling;  
- Also relevant in suspensions of microorganisms (viscous Stokes regime).
Bioturbation

The earliest case studied of animals ‘stirring’ their environment is the subject of Darwin’s last book.

This was suggested by his uncle and future father-in-law Josiah Wedgwood II, son of the famous potter.

“I was thus led to conclude that all the vegetable mould over the whole country has passed many times through, and will again pass many times through, the intestinal canals of worms.”
Munk’s Idea

Though it had been mentioned earlier, the first to seriously consider the role of ocean biomixing was Walter Munk (1966):

“...I have attempted, without much success, to interpret [the eddy diffusivity] from a variety of viewpoints: from mixing along the ocean boundaries, from thermodynamic and biological processes, and from internal tides.”
Ocean biomixing: Basic observations

The idea lay dormant for almost 40 years; then

- Huntley & Zhou (2004) analyzed swimming of 100 (!) species, ranging from bacteria to blue whales. Typical turbulent energy production is $\sim 10^{-5} \text{ W kg}^{-1}$. Total is comparable to energy dissipation by major storms.

- Another estimate comes from the solar energy captured: 63 TeraW, something like 1% of which ends up as mechanical energy (Dewar et al., 2006).

- Kunze et al. (2006) find that turbulence levels during the day in an inlet were 2 to 3 orders of magnitude greater than at night, due to swimming krill.

- However, Kunze has failed to find this effect again on subsequent cruises. Visser (2007) has questioned whether small-scale turbulence can lead to overturning.
In situ experiments

Katija & Dabiri (2009) looked at jellyfish:

[movie 1] (Palau’s Jellyfish Lake.) Correct length scale is path length?
Displacement by a moving body

Maxwell (1869); Darwin (1953); Eames et al. (1994)
A sequence of kicks

Inspired by Einstein’s theory of diffusion (Einstein, 1956): a test particle initially at $x(0) = 0$ undergoes $N$ encounters with an axially-symmetric swimming body:

$$x(t) = \sum_{k=1}^{N} \Delta_L(a_k, b_k) \hat{r}_k$$

$\Delta_L(a, b)$ is the displacement, $a_k$, $b_k$ are impact parameters, and $\hat{r}_k$ is a direction vector.

$(a > 0$, but $b$ can have either sign.)
After squaring and averaging, assuming isotropy:

\[ \langle |x|^2 \rangle = N \langle \Delta_L^2(a, b) \rangle \]

where \( a \) and \( b \) are treated as random variables with densities

\[ \frac{dA}{V} = 2 \frac{da \, db}{V} \quad (2D) \quad \text{or} \quad 2\pi a \, da \, db / V \quad (3D) \]

Replace average by integral:

\[ \langle |x|^2 \rangle = \frac{N}{V} \int \Delta_L^2(a, b) \, dA \]

Writing \( n = 1/V \) for the **number density** (there is only one swimmer) and \( N = Ut/L \) (\( L/U \) is the **time between steps**):

\[ \langle |x|^2 \rangle = \frac{Unt}{L} \int \Delta_L^2(a, b) \, dA \]
Putting this together,

\[
\langle |\mathbf{x}|^2 \rangle = \frac{2U nt}{L} \int \Delta_L^2(a, b) \, da \, db = 4\kappa t, \quad \text{2D}
\]

\[
\langle |\mathbf{x}|^2 \rangle = \frac{2\pi U nt}{L} \int \Delta_L^2(a, b) a \, da \, db = 6\kappa t, \quad \text{3D}
\]

which defines the effective diffusivity \( \kappa \).

If the number density is low \( nL^d \ll 1 \), then encounters are rare and we can use this formula for a collection of particles.
Inviscid cylinders and spheres

\[ \kappa = \frac{\pi}{3} \text{Un} \int a^2 \Delta_L^2(a, b) \, d(\log a) \, d(b/L) \]

Notice \( \Delta_L(a, b) \) is nonzero for \( 0 < b < L \); otherwise independent of \( b \) and \( L \) \( \implies \) have to cross point of closest approach.

\[ a \Delta_L^2(a, b) \text{ (cylinder)} \]

\[ a^2 \Delta_L^2(a, b) \text{ (sphere)} \]
Numerical simulation

• Validate theory using simple simple simulations;
• Large periodic box;
• $N_{\text{swim}}$ swimmers (cylinders of radius 1), initially at random positions, swimming in random direction with constant speed $U = 1$;
• Target particle initially at origin advected by the swimmers;
• Since dilute, superimpose velocities;
• Integrate for some time, compute $|\mathbf{x}(t)|^2$, repeat for a large number $N_{\text{real}}$ of realizations, and average.
A ‘gas’ of swimmers

[movie 2] 100 cylinders, box size = 1000
How well does the dilute theory work?

\[ \langle |x|^2 \rangle / 2nU^3 \]

- \( n = 10^{-3} \)
- \( n = 5 \times 10^{-4} \)
- \( n = 10^{-4} \)

\(-\) theory
Cloud of particles

[movie 3]  (30 cylinders)
Cloud dispersion proceeds by steps

\[ \langle |x|^2 \rangle \]

\( N = 30 \)
\( n = 7.5 \times 10^{-4} \)
Considerable literature on transport due to microorganisms: Wu & Libchaber (2000); Hernandez-Ortiz et al. (2006); Saintillian & Shelley (2007); Ishikawa & Pedley (2007); Underhill et al. (2008); Ishikawa (2009); Leptos et al. (2009)

Lighthill (1952), Blake (1971), and more recently Ishikawa et al. (2006) have considered squirmers:

- Sphere in Stokes flow;
- Steady velocity specified at surface, to mimic cilia;
- Steady swimming condition imposed (no net force on fluid).

(Drescher et al., 2009) (Ishikawa et al., 2006)
Typical squirmer

3D axisymmetric streamfunction for a typical squirmer, in cylindrical coordinates \((\rho, z)\):

\[
\psi = -\frac{1}{2}\rho^2 + \frac{1}{2r^3} \rho^2 + \frac{3\beta}{4r^3} \rho^2 z \left( \frac{1}{r^2} - 1 \right)
\]

where \(r = \sqrt{\rho^2 + z^2}\), \(U = 1\), radius of squirmer = 1.

\(\beta\) is the amplitude of the stresslet (distinguishes pushers/pullers).

We will use \(\beta = 5\) for most of the remainder.
Particle motion for squirmer

A particle near the squirmer’s swimming axis initially (blue) moves towards the squirmer.

After the squirmer has passed the particle follows in the squirmer’s wake.

(The squirmer moves from bottom to top.)

[movie 4]
Squirmer displacements $a^2 \Delta_{L}^2(a, b)$
Squirmers: Trajectories

The two peaks in the displacement plot come from ‘incomplete’ trajectories:

\[
\begin{align*}
&b/L = 0 \\
&b/L = 0.5 \\
&b/L = 1
\end{align*}
\]

For long path length, the effective diffusivity is independent of the swimming path length, and yet the dominant contribution arises from the finiteness of the path (uncorrelated turning directions).
Squirmers: Transport

\[ \langle |x|^2 \rangle \]

\[ x \times 10^{-4} \]

\[ t \]

\[ 0 \]

\[ 50 \]

\[ 100 \]

\[ 150 \]

\[ 200 \]
Non-Gaussian PDFs of displacement

- Variance exhibits similar short-time ballistic scaling as in Wu & Libchaber (2000) (due to smoothness);
- PDF qualitatively matches experiments of Leptos et al. (2009). In our case, exponential tails are due to sticking at the stagnation points on the squirmer’s body.
Conclusions

- Simple **dilute model** works well for a range of swimmers;
- Slip surfaces have an effective diffusivity that is **independent of path length**, for long path length;
- No-slip flows dominated by **sticking** and have a **log dependence** on path length;

Future work:
- Wake models and turbulence;
- PDF of scalar concentration;
- **Buoyancy effects** for the ocean case;
- Higher densities;
- Schooling: longer length scale?


